

Electronic Transport Phenomena in Composite Nanocrystalline/Amorphous and Free-Standing Nanocrystalline Thin Films

James Kakalios

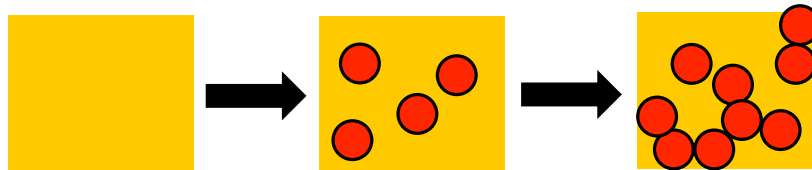
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Acknowledgements

- All of this research performed in collaboration with **Prof. Uwe R. Kortshagen**,
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NINN Characterization Facility, **Minnesota Nanofabrication
Center** and **The University of Minnesota**

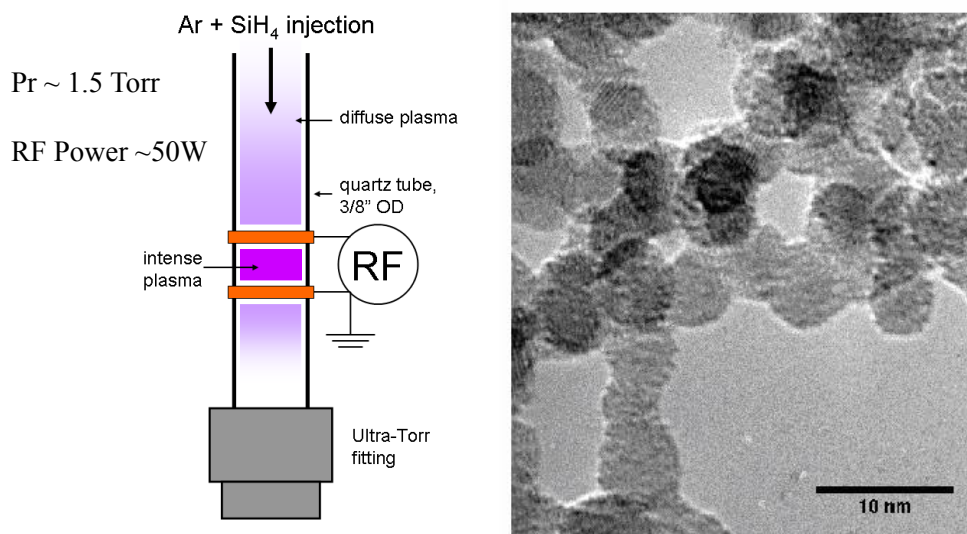
Composite Amorphous/ Nanocrystalline Thin Films

- Thin Film Amorphous Matrix in which Nanocrystallites are Embedded ($X_c \sim 0 - 100\%$)
- Composites Can Combine the Large Area, Thin Film Advantages of Amorphous Silicon with Superior Electronic Properties of Crystalline Silicon
- How Do Transport Properties Change with nc?



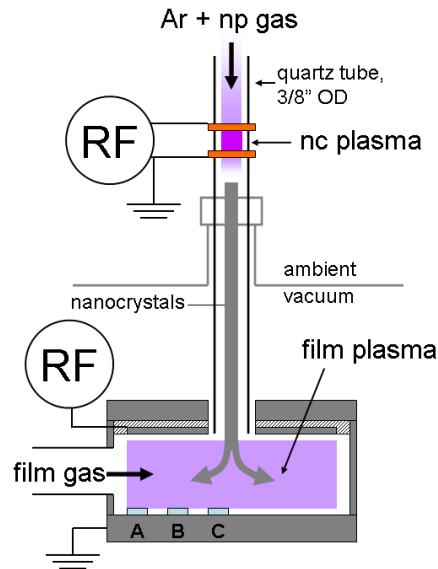
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PECVD Synthesis of Nanocrystals



Dual Chamber Co-Deposition System

- Nanocrystals are produced in one capacitively coupled PECVD chamber
- In the second chamber a-Si:H film is produced
- The nanoparticles are entrained by a carrier gas and injected in the second chamber where mixed phase film is deposited
- Growth parameters can be separately optimized for both the particles and the amorphous film

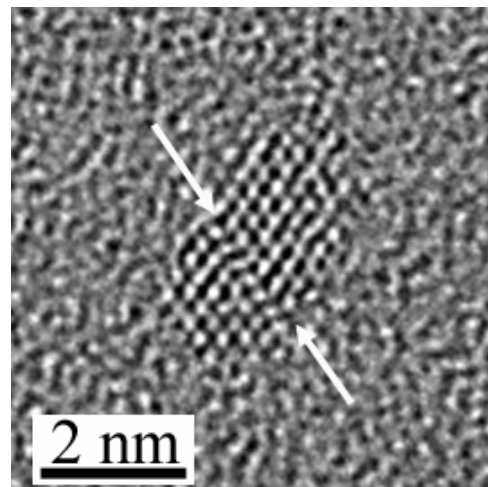


High Resolution TEM

Si particle $\sim 1.5 \text{ nm} \times 3 \text{ nm}$

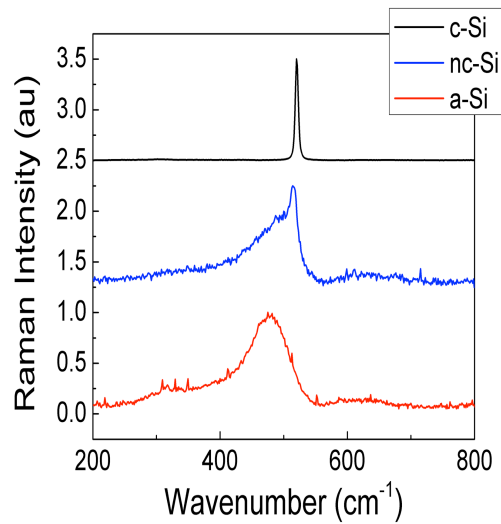
Images taken by Prof. C. B. Carter
with a Philips CM 200 FEG
with a spherical aberration
corrector

with Dr. Markus Lentzen and Prof.
Knut Urban (Research Center
Jülich, Germany).

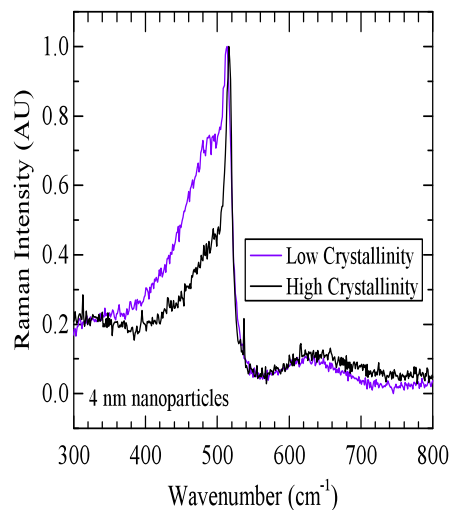


HRTEM image of the
1450 mTorr a/nc-Si:H sample

Raman Spectrum for c-Si, nc-Si and a-Si

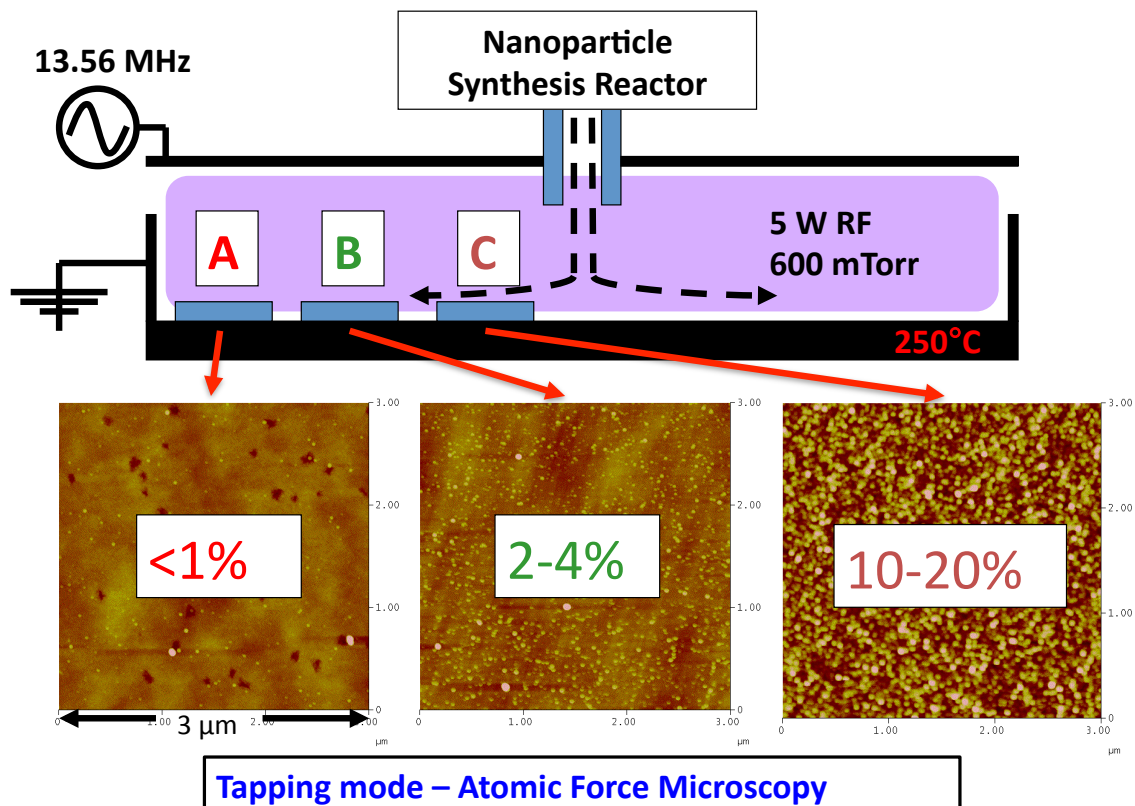
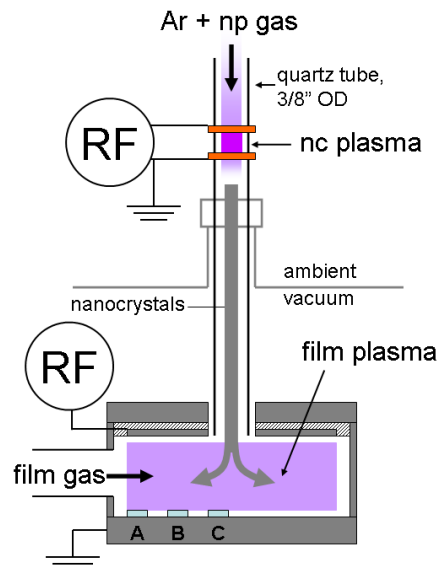


Raman for nc-Si with varying degree of crystallinity



Nanocrystal Concentration is **Not** Uniform

- Concentration of nc in a-Si:H film controlled through gas convection in second chamber

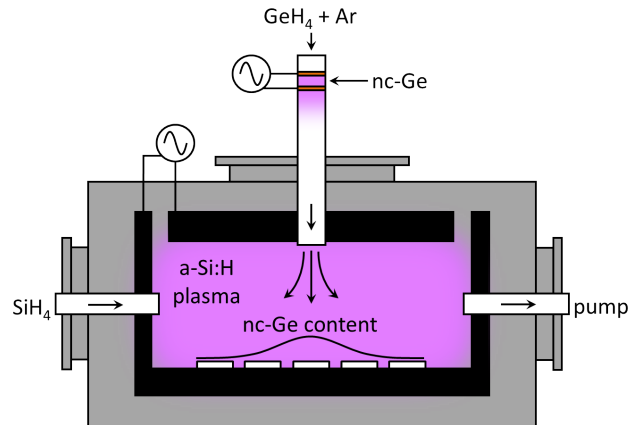
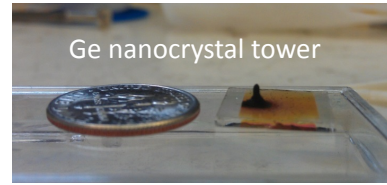


Germanium nc in a-Si:H

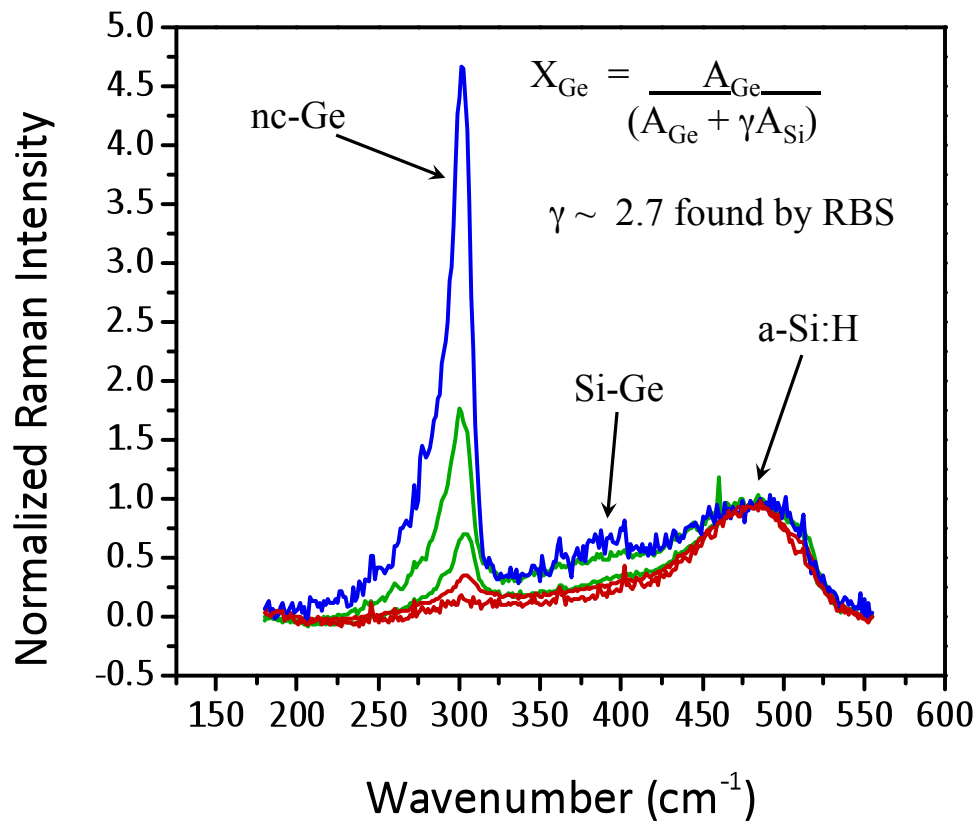
~4nm Ge crystals embedded in a-Si:H matrix

Vary only nc-Ge content across a series of films

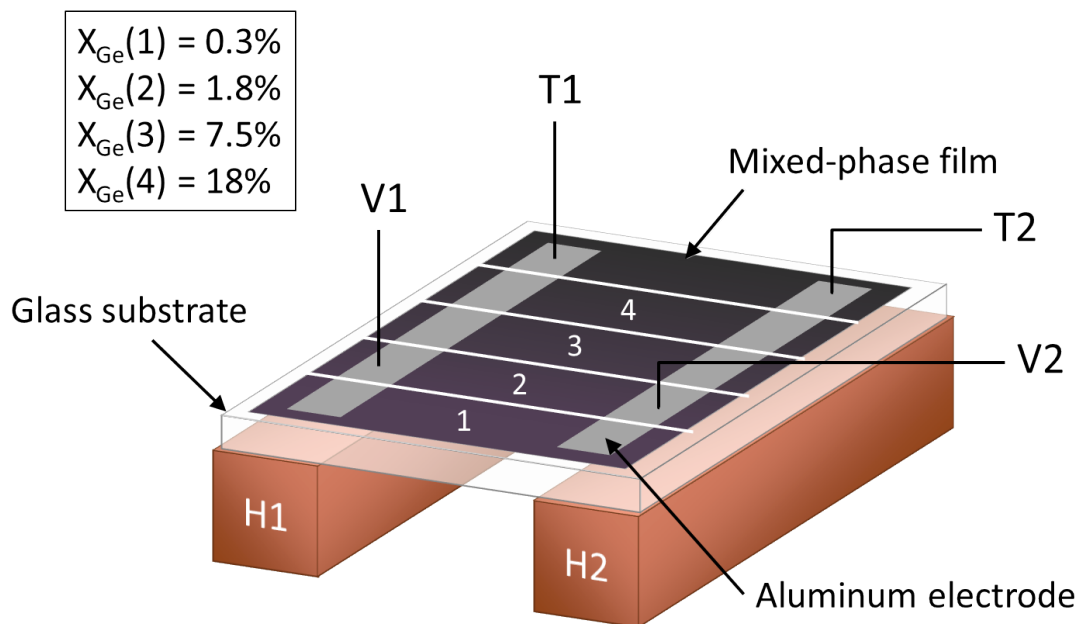
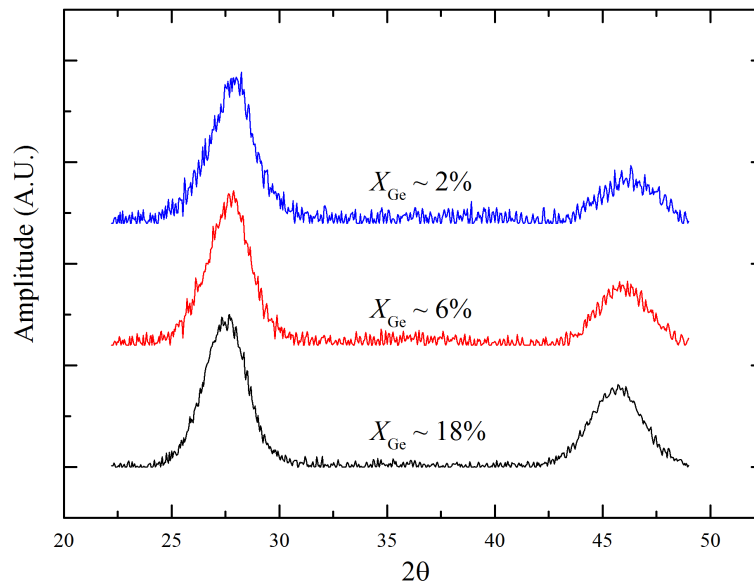
Over multiple runs, nc-Ge content varies from 0 to 75%

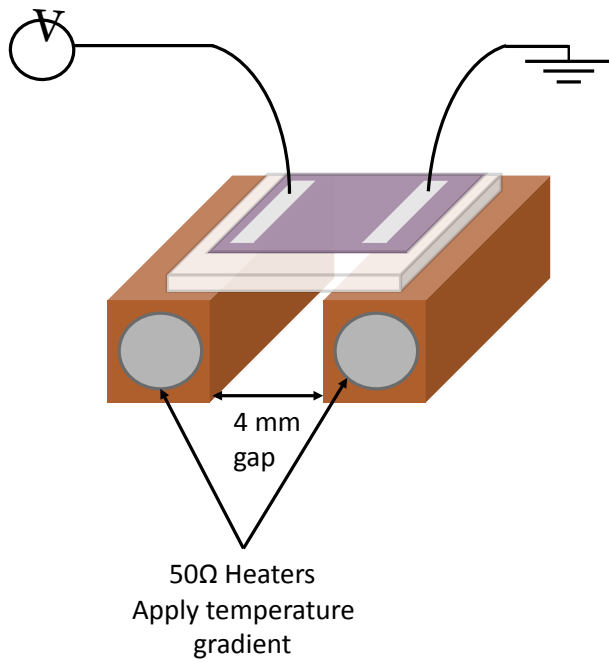


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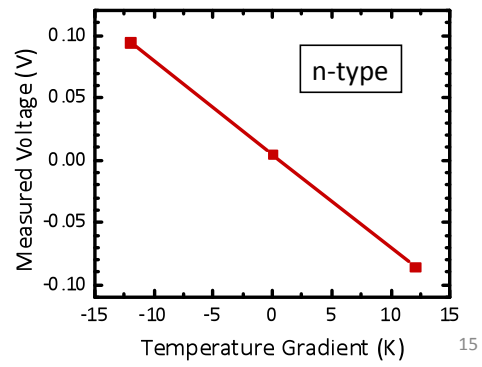
X-Ray Diffraction Indicates nc diameter $\sim 3.5 - 4.5$ nm



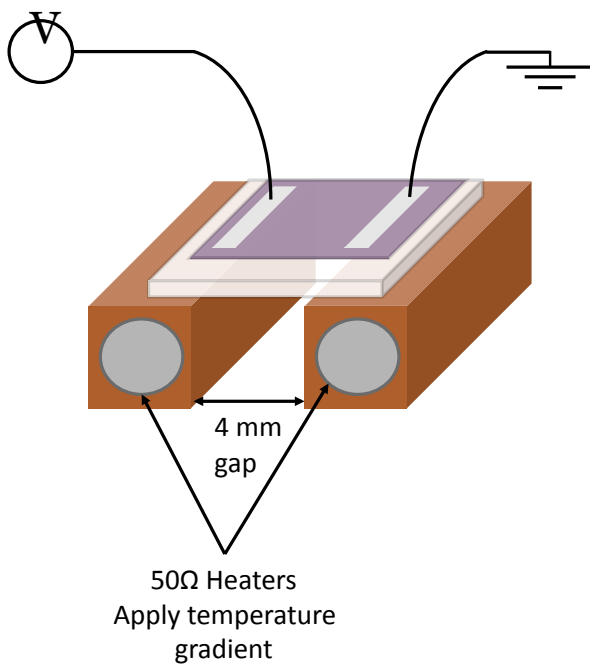


$$\text{Seebeck coefficient } S = -\frac{V_1 - V_2}{T_1 - T_2}$$

Step	T_1	T_2
1)	$T_0 + 6$	$T_0 - 6$
2)	T_0	T_0
3)	$T_0 - 6$	$T_0 + 6$

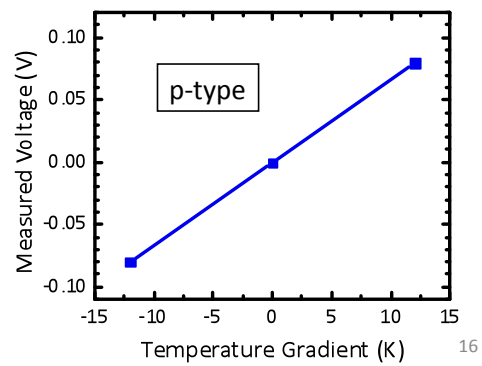


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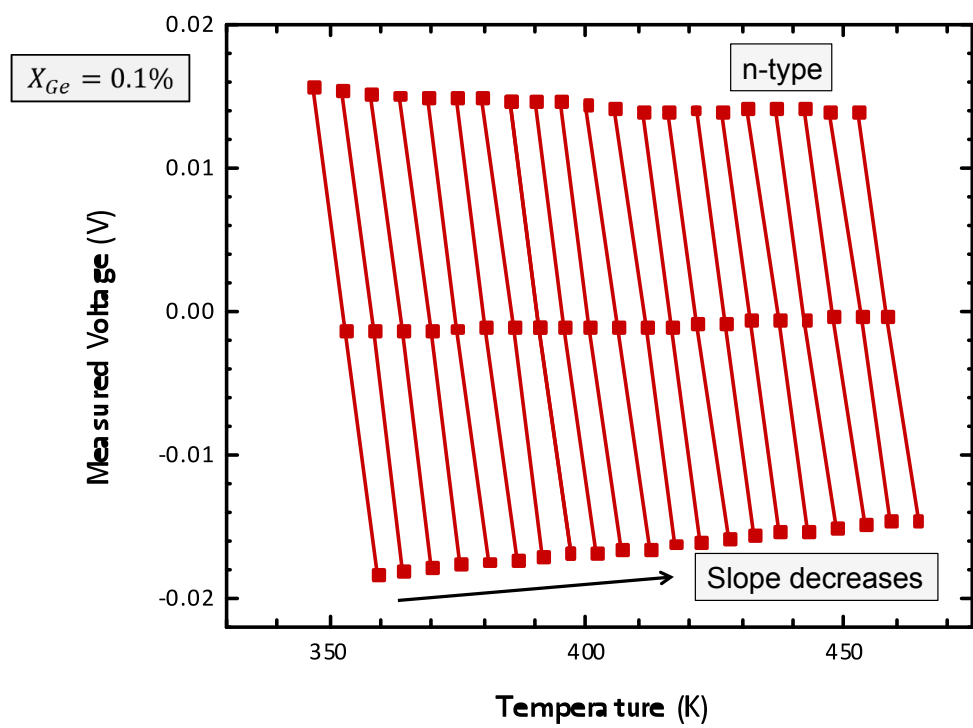


$$\text{Seebeck coefficient } S = -\frac{V_1 - V_2}{T_1 - T_2}$$

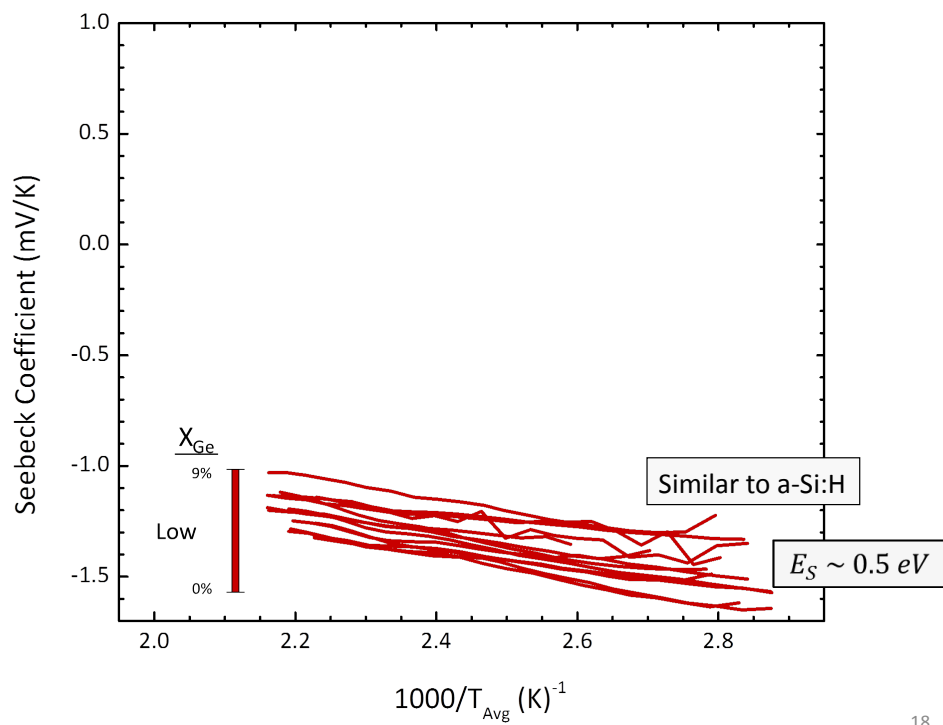
Step	T_1	T_2
1)	$T_0 + 6$	$T_0 - 6$
2)	T_0	T_0
3)	$T_0 - 6$	$T_0 + 6$



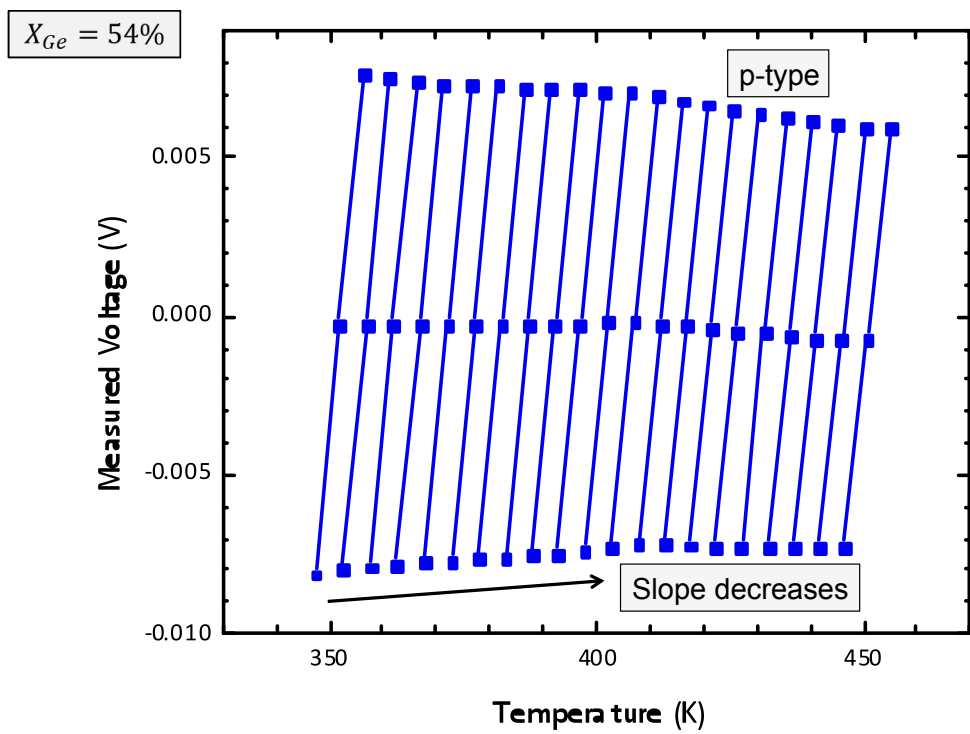
16



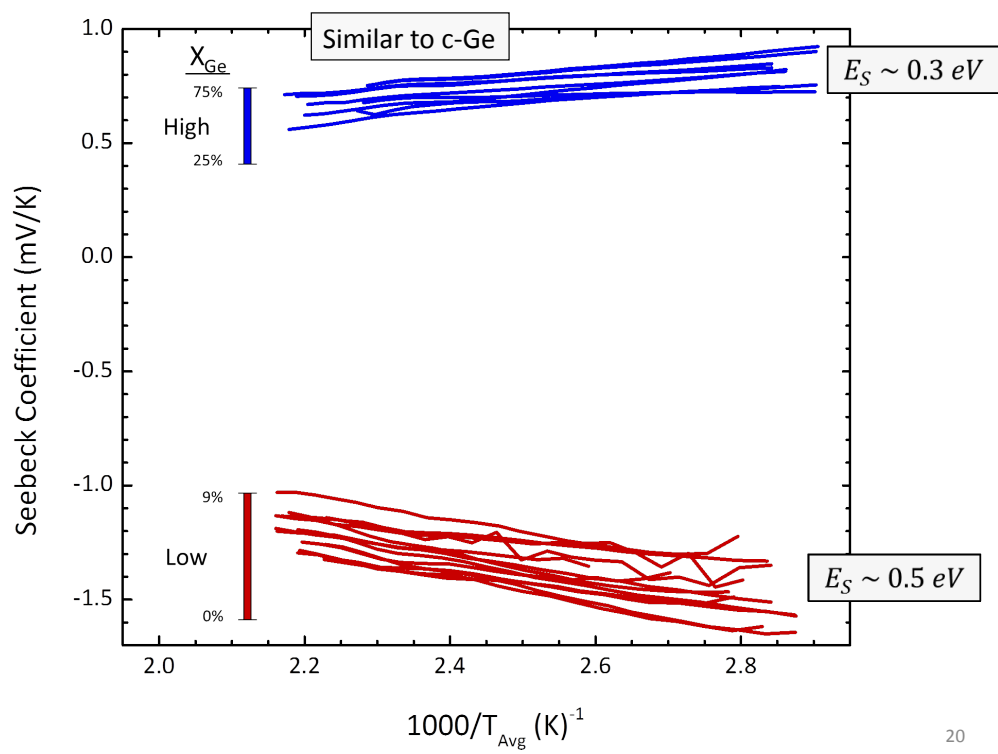
17



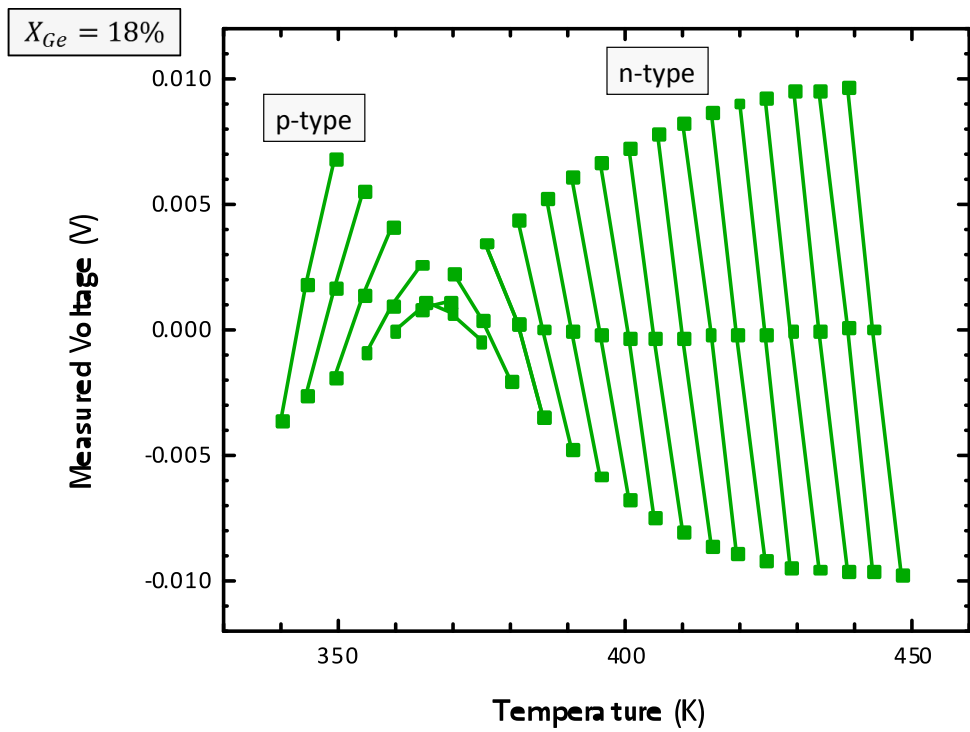
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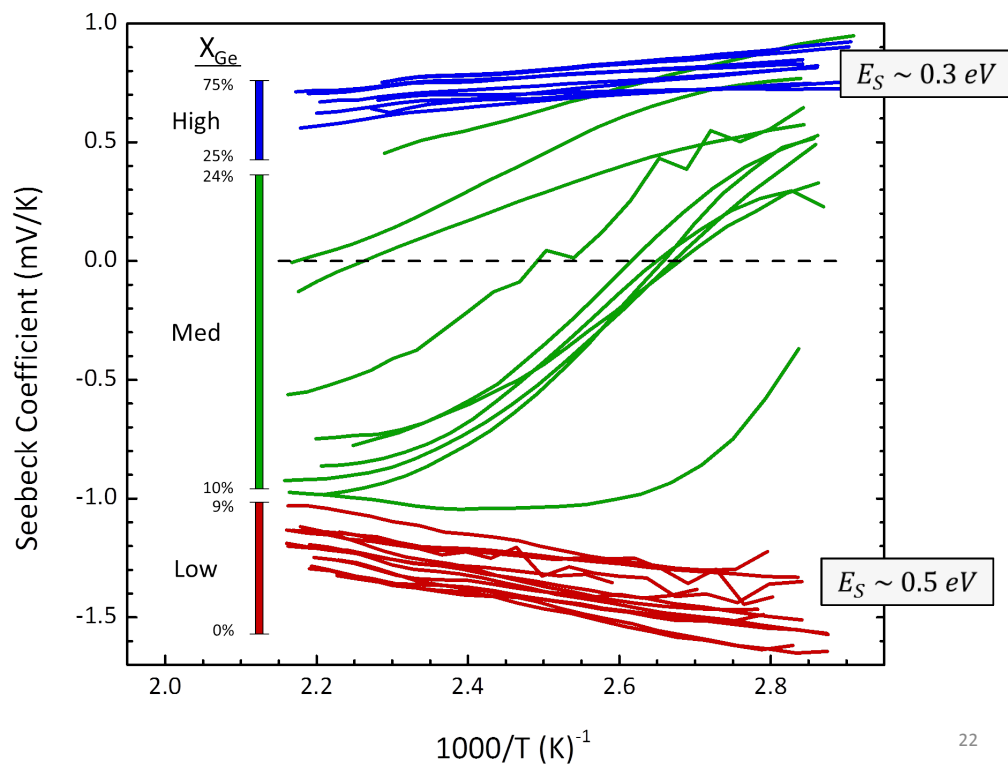


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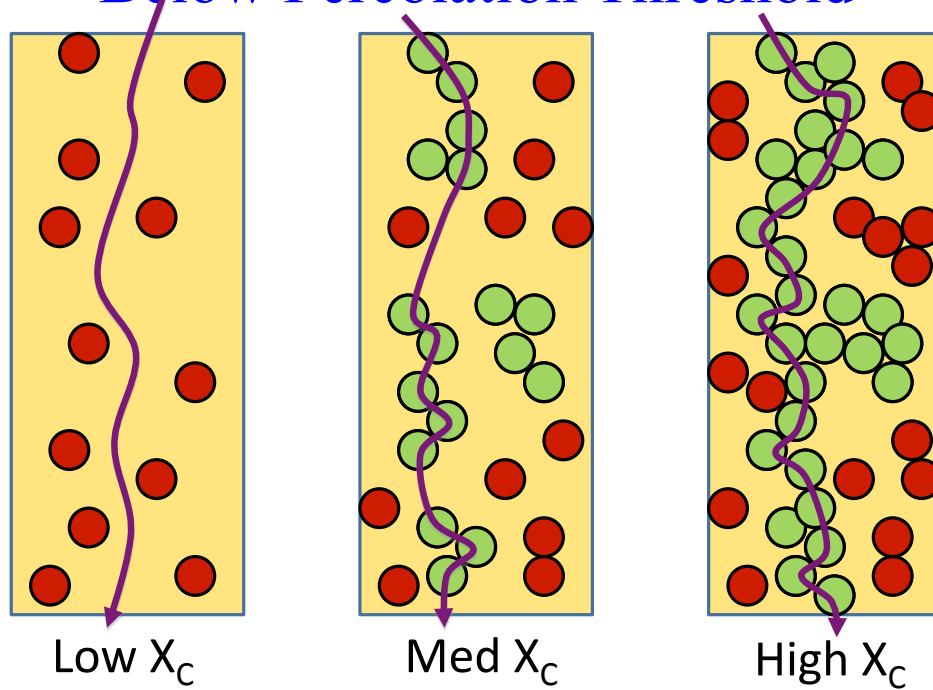
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K. Bodurtha and JK, J. Appl. Phys. **114**, 193705 (2013)



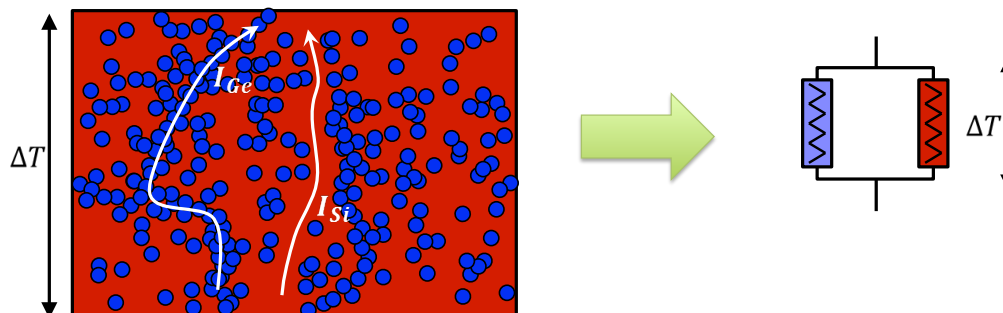
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Can See Influence of nc-Ge Chains Below Percolation Threshold



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Parallel Resistor Model



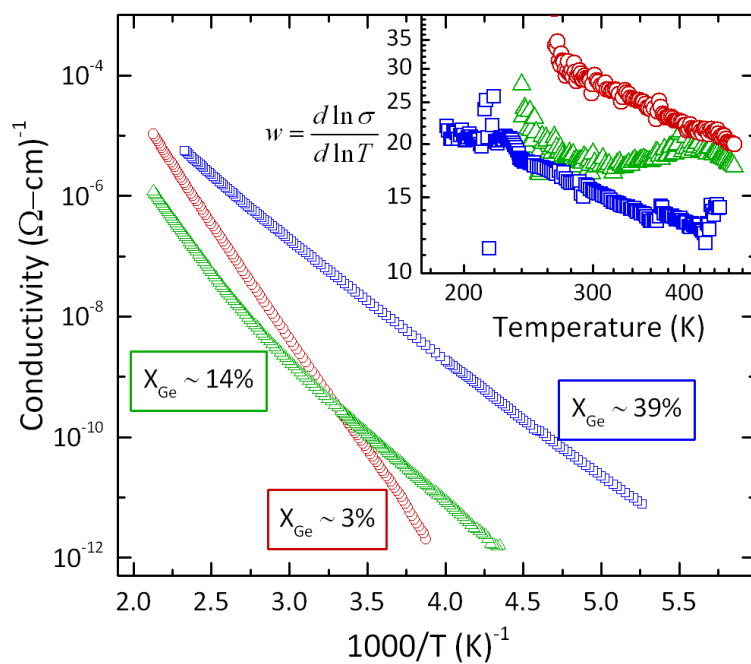
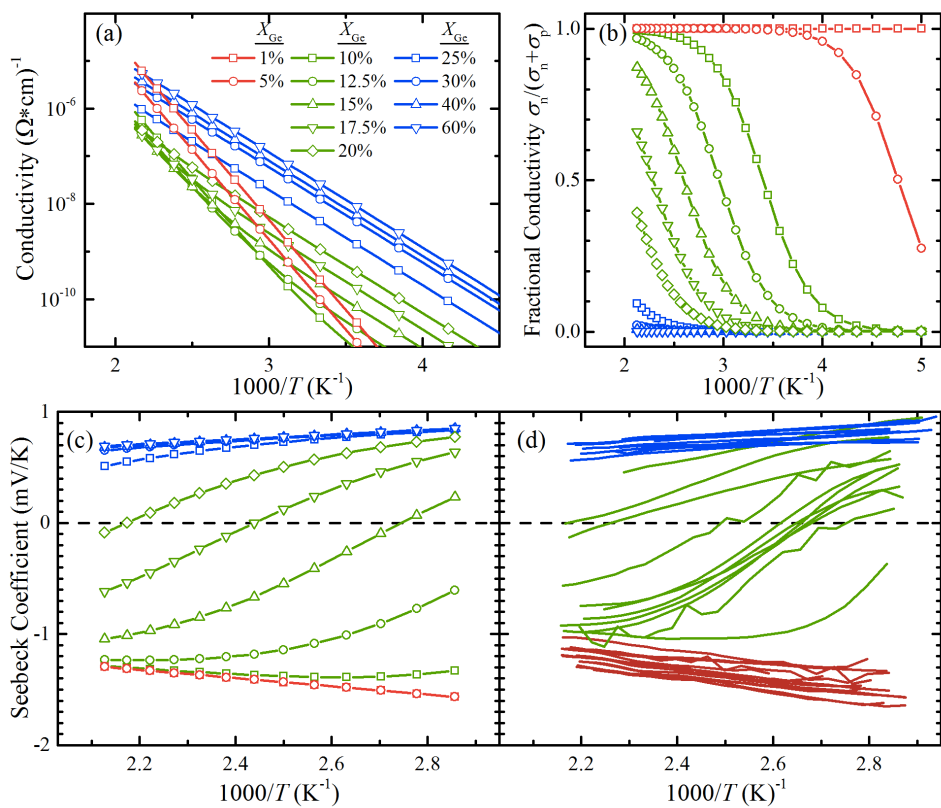
Conductivity is weighted by the cross-sectional area of each resistor, which is proportional to X_{Ge}

$$\sigma(T) = X_{Ge}\sigma_{Ge} + (1 - X_{Ge})\sigma_{Si}$$

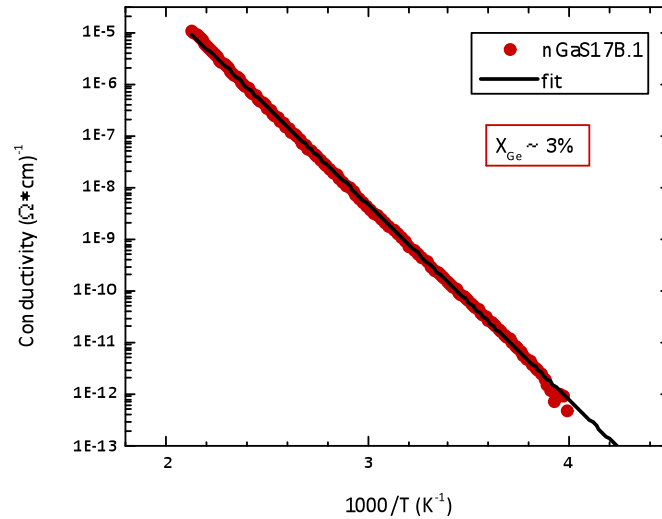
Thermopower is weighted by the conductivity:

$$S(T) = \frac{X_{Ge}\sigma_{Ge}S_{Ge} + (1 - X_{Ge})\sigma_{Si}S_{Si}}{\sigma(T)}$$

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Is Conductivity of nc-Ge/a-Si:H Arrhenius?

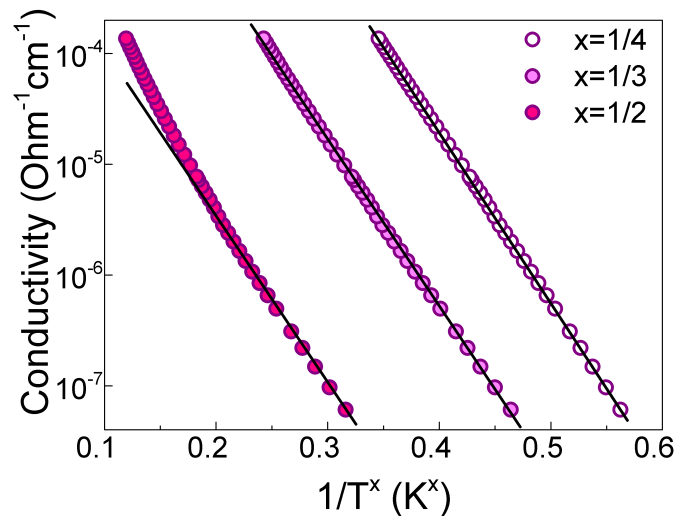


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How to Determine Temperature Dependence of Conductivity?

Change plot format to yield straight line

Often ambiguous - Temperature range too small



How to Determine Temperature Dependence of Conductivity?

Reduced Activation Energy Analysis

Zabrodskii and Shlimak, *Sov. Phys. Semicond.* (1975); **Hill**, *Phys. Stat. Solidi* (1976).

$$\sigma = \sigma_0 \exp [- (T_0 / T)^x]$$

Calculate $w(T)$ from $\sigma(T)$ data

$$w(T) = \frac{d \ln(\sigma)}{d \ln(T)} \quad \longrightarrow \quad w(T) = x \left(\frac{T_0}{T} \right)^x$$

Reduced Activation Energy Reflects Conduction Mechanism

$$\sigma = \sigma_0 \exp [- (T_0 / T)^x] \quad \longrightarrow \quad w(T) = x \left(\frac{T_0}{T} \right)^x$$

Simple Activated Behavior $x = 1$ $\log W(T) \sim (-1) \log T$

Mott – Variable Range Hopping $x = 1/4$ $\log w(T) \sim (-1/4) \log T$

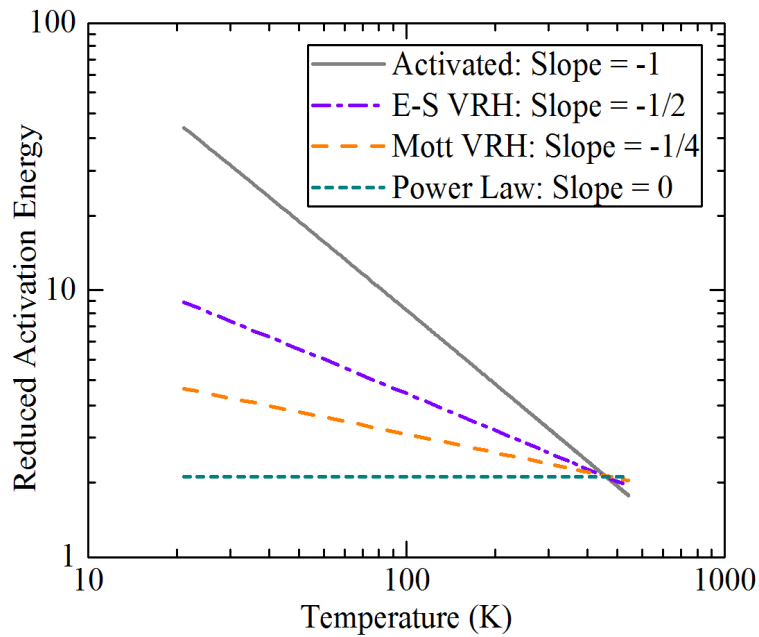
Efros-Shklovskii VRH $x = 1/2$ $\log w(T) \sim (-1/2) \log T$

Power Law Temp Dep. $\sigma \sim T^n$ $d \log \sigma / d \log T = W = n$

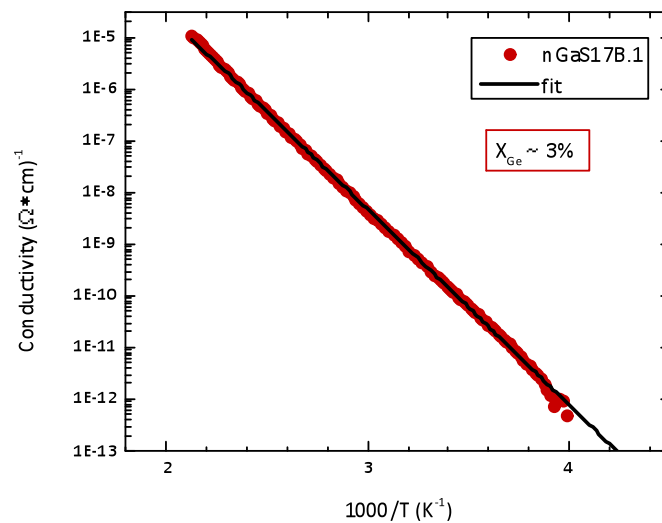
Zabrodskii and Shlimak, *Sov. Phys. Semicond.* (1975); **Hill**, *Phys. Stat. Solidi* (1976).

Log-Log Plot of Reduced Activation Energy against Temperature Indicates Power-Law Exponent

$$w(T) = d\ln(\sigma)/d\ln(T)$$

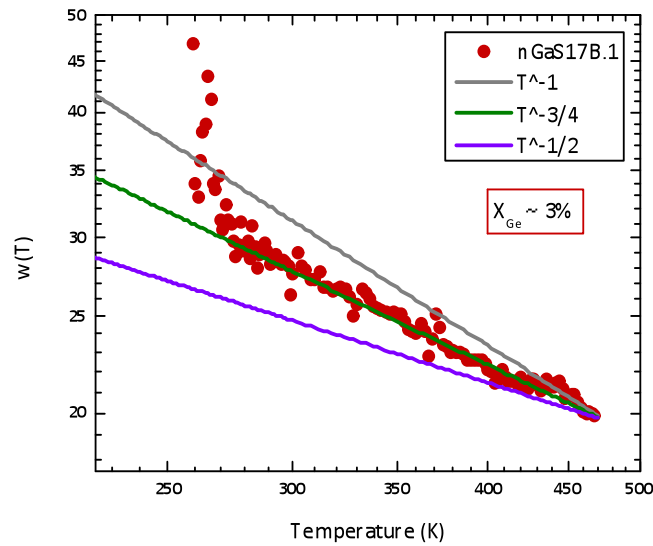


Apply the Zabrodskii Analysis Technique to nc-Ge/a-Si:H



Zabrodskii Analysis Finds Best Fit for $\kappa \sim 3/4$

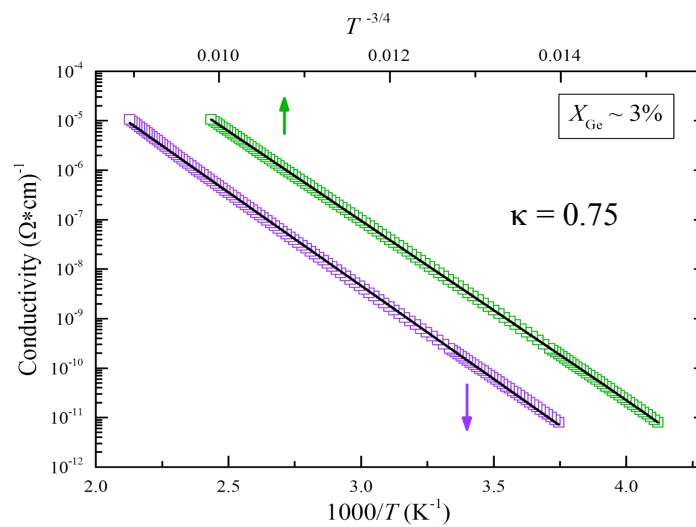
$$w(T) = d \ln(\sigma) / d \ln(T)$$



K. Bodurtha and JK, J. Appl. Phys. **118**, 215103 (2015)

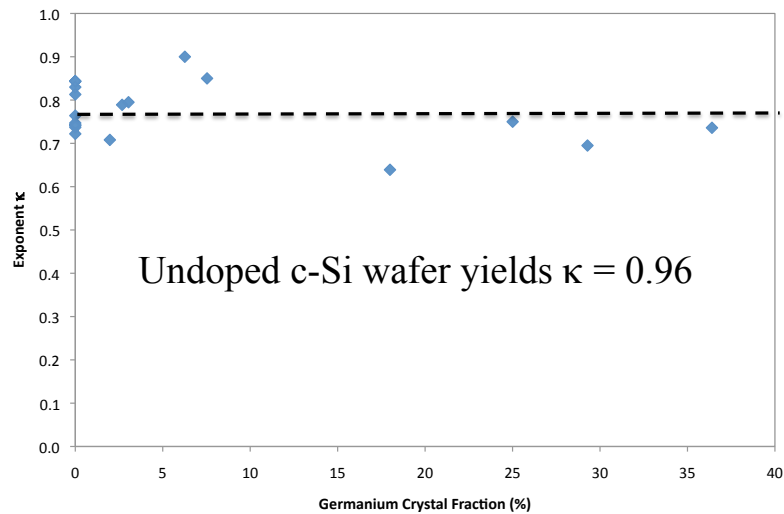
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$$\sigma = \sigma_0 \exp [-(T_0/T)^\kappa]$$

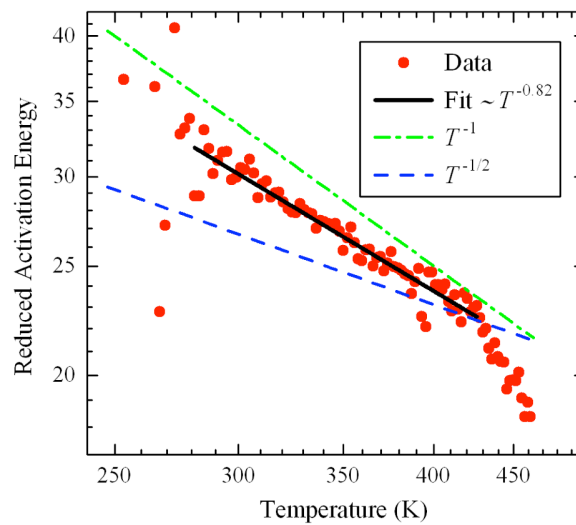


Power-Law Exponent κ vs. nc-Ge Content

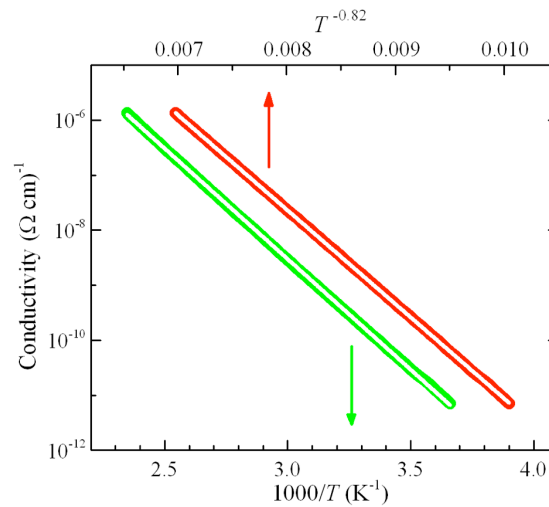
$$\sigma = \sigma_0 \exp [-(T_0/T)^\kappa]$$
$$0.7 < \kappa < 0.84$$



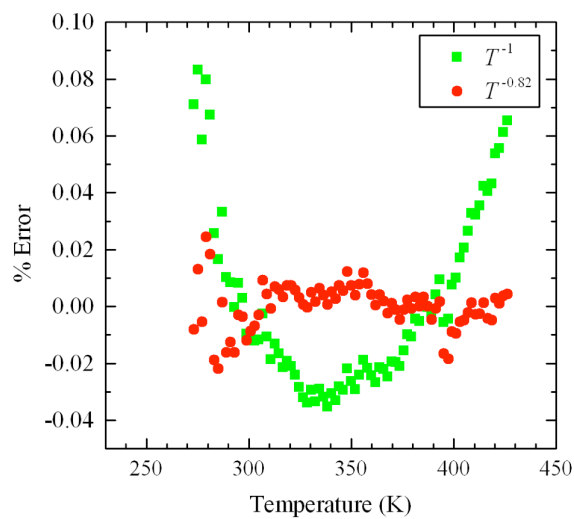
Reduced Activation Energy Plot for a-Si:H



Arrhenius vs. Anomalous Hopping Temperature Dependence



Difference between Data and Arrhenius and Anomalous Hopping Expression



Hopping in Exponential Band Tail States

Current is carried by free electrons thermally excited into the conduction band.

Bandtails don't contribute to conduction because DOS is too small and states are too far apart for wavefunctions to overlap.

$$\sigma(T) = \sigma_0 \exp\left[-\frac{E_\sigma}{k_B T}\right]$$

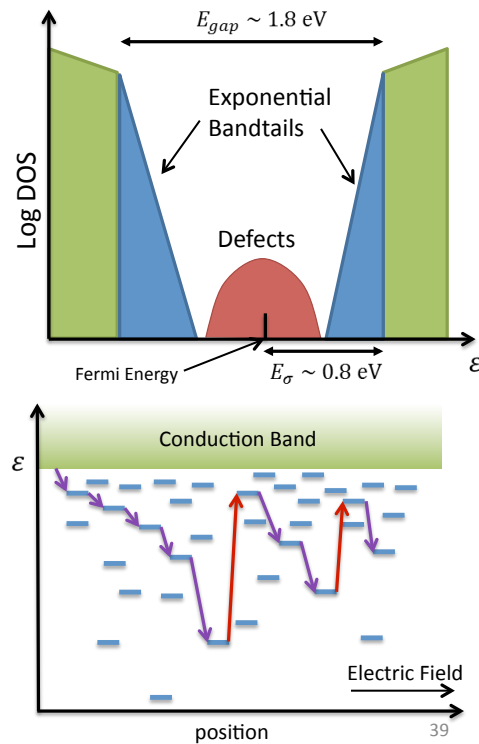
N. F. Mott, Advances in Physics 16, 49 (1967).

If bandtails are sufficiently broad, and/or T is small, hopping through bandtail states can become significant

Hop to nearby states, always hopping down in ϵ

Occasionally thermally promoted back up to higher energies.

Grunewald and Thomas (1979); D. Monroe, (1985).



Hopping Through Bandtail States

Simulations based upon model of hopping through exponential bandtails

Assume conduction occurs through band of states at transport energy of width $W = (6\epsilon_0 kT)^{1/2}$

Shklovskii, Levin, Fritzsche and Baranovskii (1990)

Hopping Through Bandtail States

Simulations based upon model of hopping through exponential bandtails

Assume conduction occurs through band of states at transport energy of width
 $W = (6\epsilon_0 kT)^{1/2}$

Shklovskii, Levin, Fritzsche and Baranovskii (1990)

For $\epsilon_0 = 50$ meV

$N = 10^{20} \text{ cm}^{-3}$

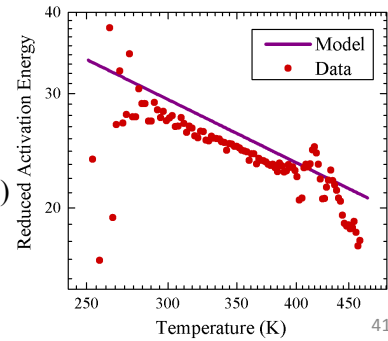
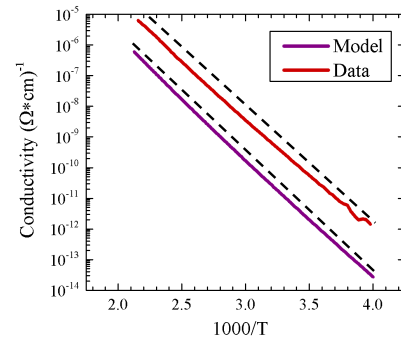
$a = 1.5 \text{ nm}$ (dist between adjacent states)

$\epsilon_F = 0.85 \text{ eV}$

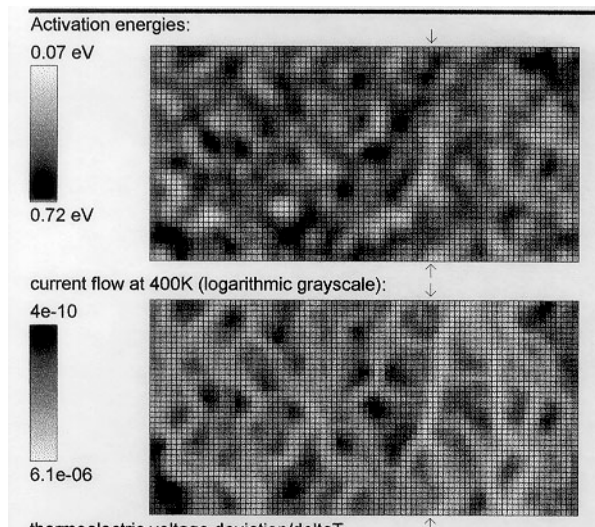
$\nu_0 = 10^{15} \text{ sec}^{-1}$

$\kappa = 0.78$ (model)

$\kappa = 0.76$ (data)



Simulations Show Current Filaments Arise from Spatial Variations of Resistance

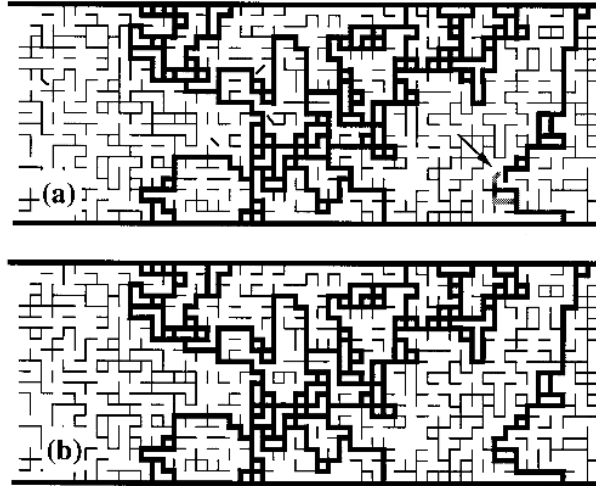


X-Y Grid of Resistors

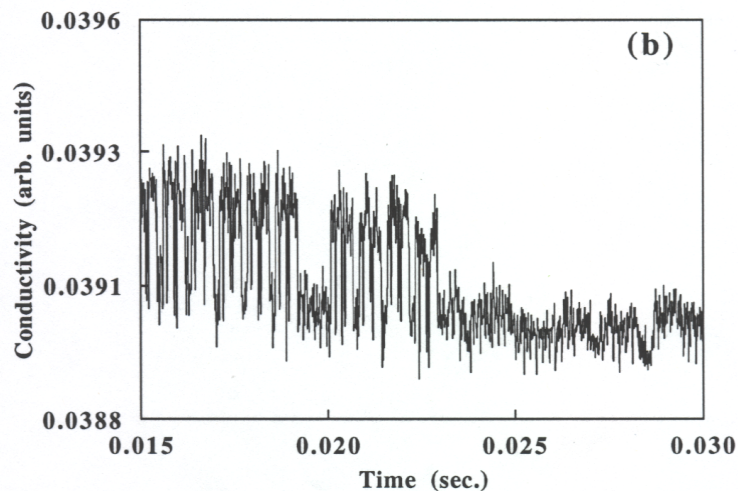
$$R = R_0 \exp[E_a/kT]$$

D. Quicker and JK,
 Phys Rev B **60** (1999)

Dynamical Percolation Model Simulates Effect of H Motion on Inhomogeneous Current Filaments

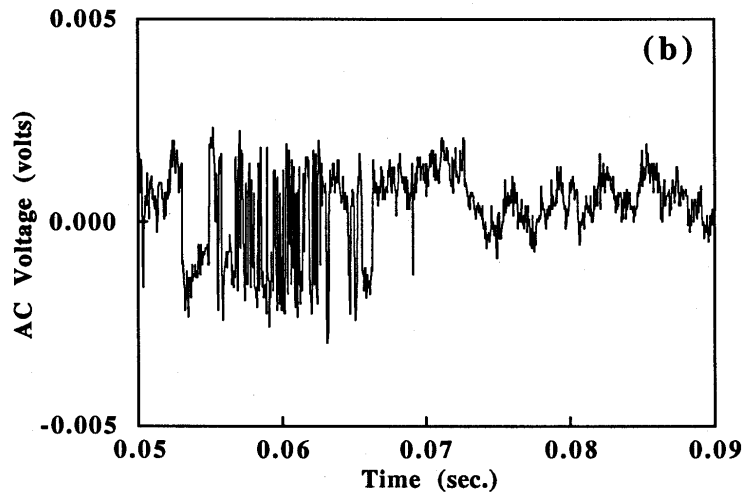


Simulated Current Fluctuations Show Both RTSN and $1/f$ Noise

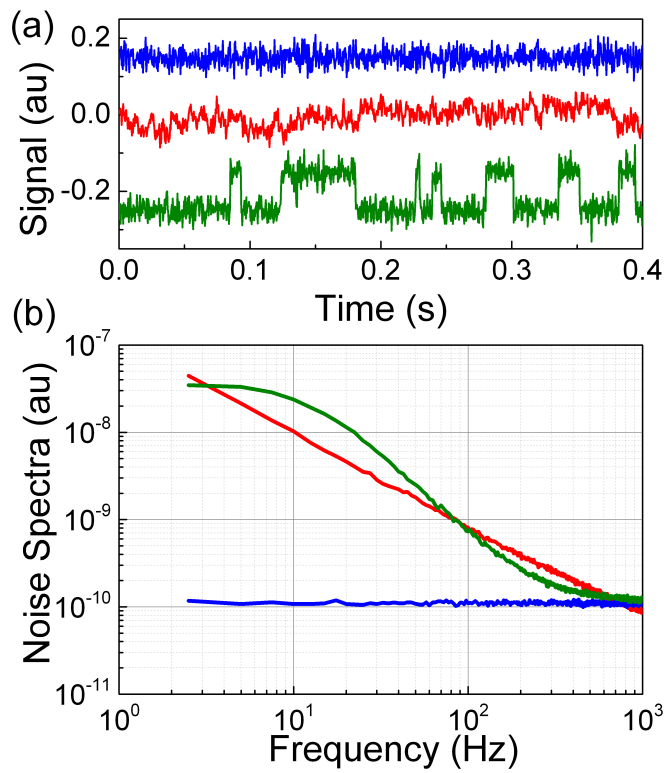


L. Lust and JK, Phys Rev E (1994); Phys Rev Lett **75** (1995)

Consistent with Measured Current Fluctuations



L. Lust and JK, Phys Rev E (1994); Phys Rev Lett **75** (1995)



Correlation Coefficients Quantify Interactions Across Frequency Octaves

$$\rho_{ij} = \frac{\sum (NP_i - \langle NP_i \rangle)(NP_j - \langle NP_j \rangle)}{(K - 1) \sigma_i \sigma_j}$$

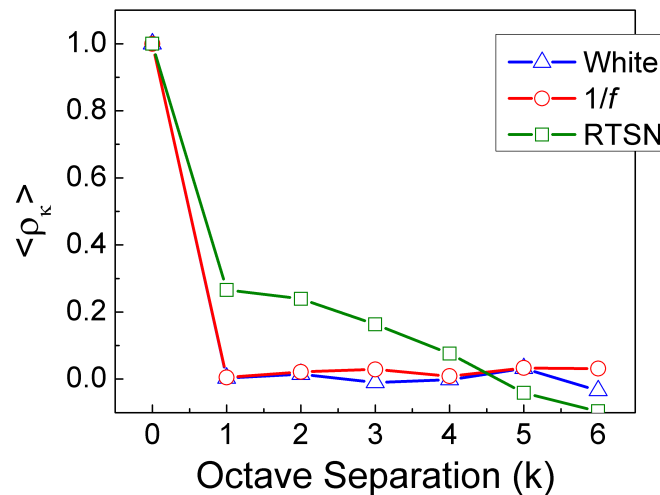
NP_i = Noise Power in Octave i ($i = 1 - 7$)

$\langle NP_i \rangle$ = Average Noise Power in Octave i

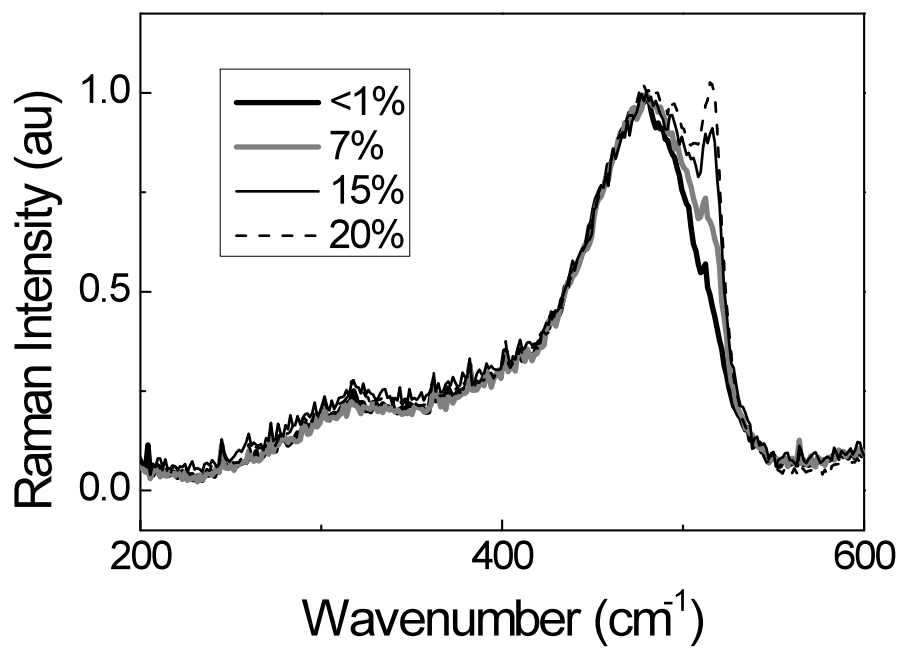
σ_i = Standard Deviation of Average Noise Power in Octave i

$K = 1 - 1024$ FFT's

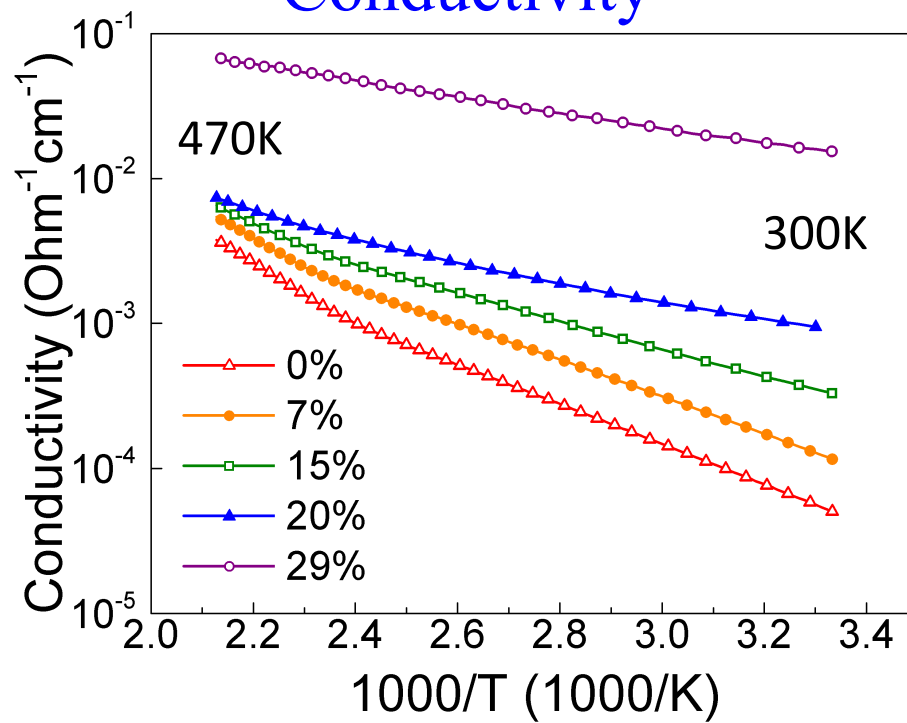
Simulated Correlation Coefficients

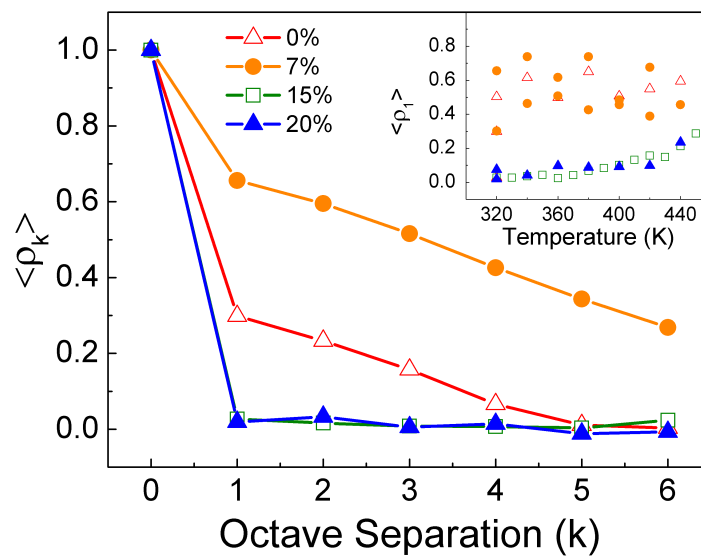
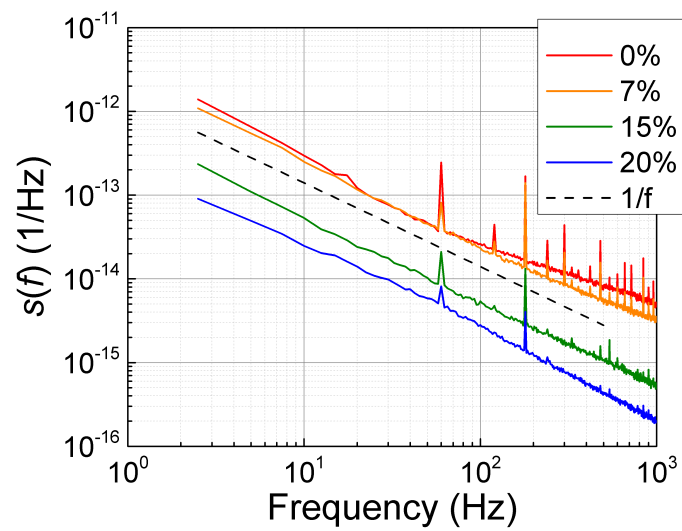


N-Type Doped a/nc-Si:H Films – Raman Spectra

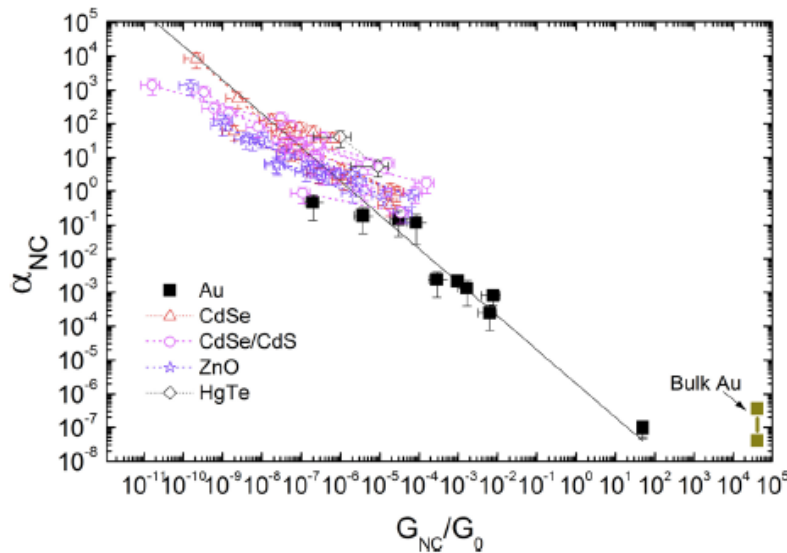


Conductivity





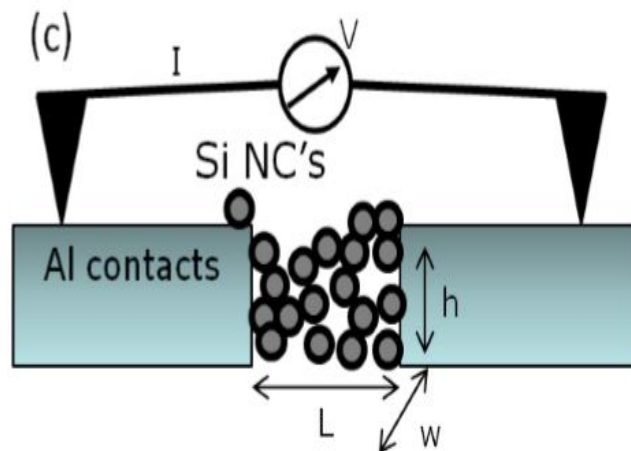
$1/f$ Noise in QD Very Sensitive to Nearest Neighbor Conductance



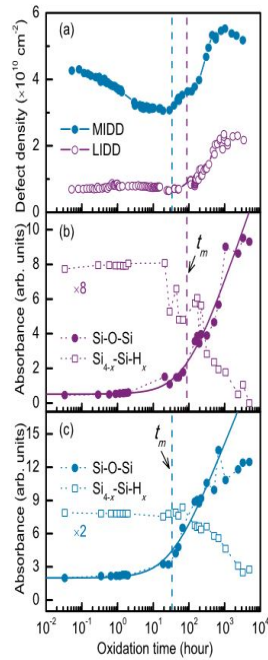
H. Liu, E. Lhuillier, and P. Guyot-Sionnest, J. Appl. Phys. **115**, 154309 (2014)

Free-Standing Nanocrystal Films

- Substrate with Pre-Deposited Electrodes Underneath Injection Tube from Nanocrystal PECVD
- Si Nanocrystals deposited on top of electrodes

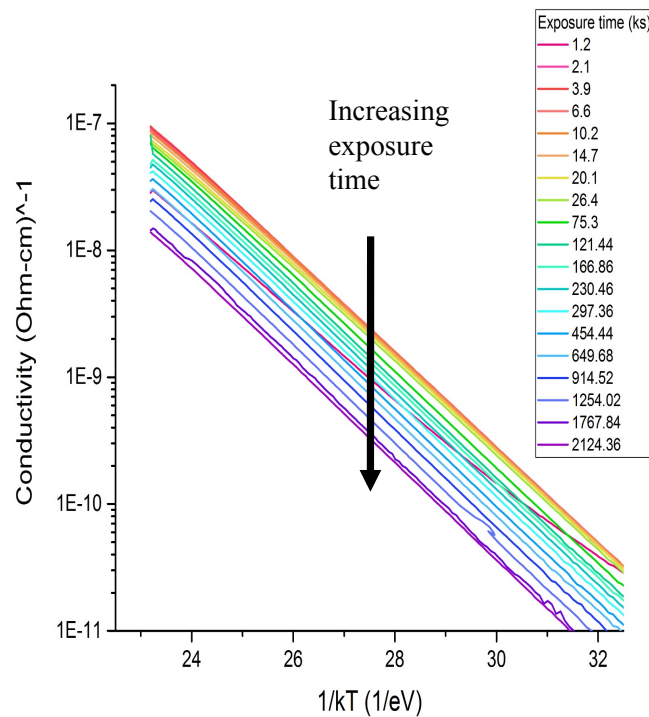


Oxides Grow Slowly on Si nc

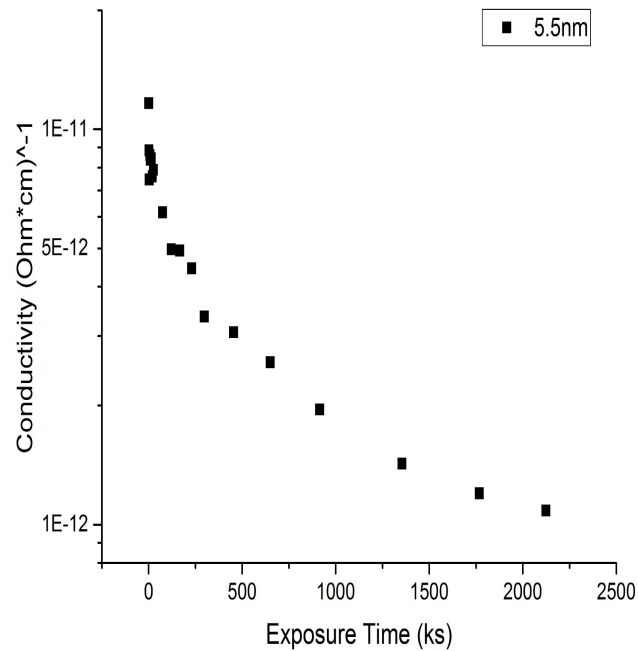


R.N. Pereira, D. J. Rowe, R. J. Anthony, and U. Kortshagen. Phys. Rev. B **86** 085449(2012).

Dark Conductivity Decreases With Atmosphere Exposure

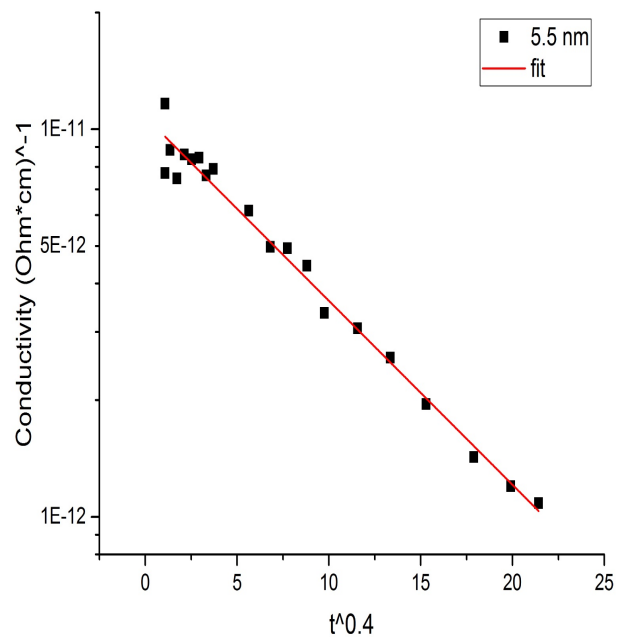


Conductivity Decreases with Increasing Air Exposure

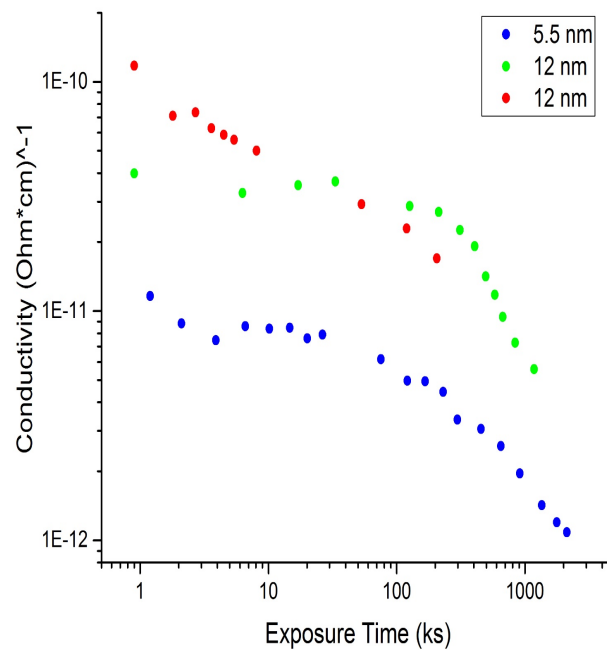


Conductivity Follows Stretched Exponential Dependence on Air Exposure Time

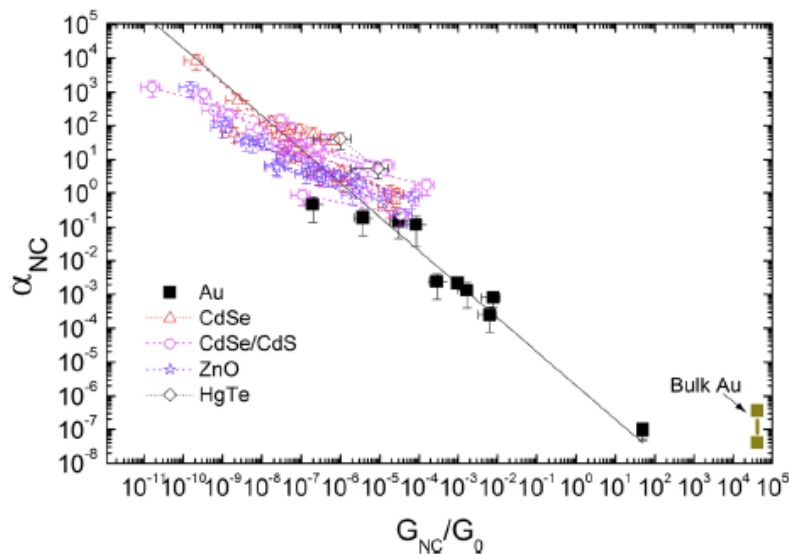
- Stretched exponential:
$$\sigma = \sigma_0 e^{-(t/\tau)^\beta}$$
- Here $\beta = 0.4$
- $\tau = 255 \text{ ks}$



Decrease of Conductivity with Atmosphere Exposure Sensitive to nc Size



1/f Noise in QD Very Sensitive to Nearest Neighbor Conductance



Summary

- Conduction In Amorphous/Nanocrystalline Composites Can Display Behavior Not Seen in Either Material Separately
- Anomalous Hopping ($\kappa \sim 3/4$) Observed in nc-Ge/a-Si:H and a-Si:H
- Transport in Free-Standing Si nc Films is Sensitive to Air Exposure